

THE ANALYSIS OF IMAGES IN N -POINT GRAVITATIONAL LENS BY METHODS OF ALGEBRAIC GEOMETRY

ALBERT T. KOTVYTSKIY^{a,b,*}, SEMEN D. BRONZA^b,
VOLODYMYR YU. SHABLENKO^a

^a Karazin Kharkov National University, Svobody Square 4, Kharkiv, 61022, Ukraine

^b Ukrainian State University of Railway Transport, Feierbakh Square 7, 61050, Kharkiv, Ukraine

* corresponding author: kotvytskiy@gmail.com

ABSTRACT. This paper is devoted to the study of images in N -point gravitational lenses by methods of algebraic geometry. In the beginning, we carefully define images in algebraic terms. Based on the definition, we show that in this model of gravitational lenses (for a point source), the dimensions of the images can be only 0 and 1. We reduce it to the fundamental problem of classical algebraic geometry - the study of solutions of a polynomial system of equations. Further, we use well-known concepts and theorems. We adapt known or prove new assertions. Sometimes, these statements have a fairly general form and can be applied to other problems of algebraic geometry. In this paper, the criterion for irreducibility of polynomials in several variables over the field of complex numbers is effectively used. In this paper, an algebraic version of the Bezout theorem and some other statements are formulated and proved. We have applied the theorems proved by us to study the imaging of dimensions 1 and 0.

KEYWORDS: gravitational lense, images, algebraic geometry, resultant.

1. INTRODUCTION

In modern astrophysics, gravitational lensing has been transformed from an effect that confirms the general theory of relativity to the research tool. Gravitational lensing is used to study both stellar systems and planets in them, and galaxies and systems of galaxies. Even the cosmological parameters of the entire metagalaxy are investigated.

From this point of view, it seems rather strange that until now a complete analytical description has been performed only for the simplest lenses - axially symmetric lenses (see for example [1] or straight infinite cosmic strings [2]).

To analyze fairly simple 2-point gravitational lenses, only approximate or numerical methods are used [3, 4].

In this paper, the authors continue the analytic study of N -point gravitational lenses by methods of algebraic geometry [5–8].

In physics, the concept of “image in a gravitational lens” is understood intuitively and is usually not determined. However, the absence of a definition can lead to ambiguous understanding of the concept and a different interpretation of some results, for example, the theorem on the odd number of images [9, 10]. On the other hand, the terminology developed in algebraic geometry makes it possible to pinpoint the concept of an image in a gravitational lens. On this basis it is possible to formulate a number of statements.

2. THE PHYSICAL FORMULATION OF THE PROBLEM FROM AN ALGEBRAIC POINT OF VIEW

For the model of a plane gravitational lens, we can write the equation that connects the coordinates of the

source (the radius vector \vec{y}) and the image coordinates (radius vector \vec{x}), see [9, 10]

$$\vec{y} = \vec{x} - \vec{\alpha}, \tag{1}$$

where $\vec{\alpha}$ is the total angle of deflection of the light beam in the plane of the lens. In the case of an N -point gravitational lens, the deflection angle is determined by the following expression:

$$\vec{\alpha} = \sum_{i=1}^N m_i \frac{\vec{x} - \vec{l}_i}{|\vec{x} - \vec{l}_i|^2}, \tag{2}$$

where m_i are dimensionless masses whose position in the plane of the lens is determined by the radius vectors \vec{l}_i . We have that holds $\sum_{i=1}^N m_i = 1$.

Equation (1) with allowance for (2) in coordinate form has the form

$$\begin{cases} \left(x_1 - \sum_{i=1}^N m_i \frac{x_1 - a_i}{(x_1 - a_i)^2 + (x_2 - b_i)^2} \right) - y_1 = 0, \\ \left(x_2 - \sum_{i=1}^N m_i \frac{x_2 - b_i}{(x_1 - a_i)^2 + (x_2 - b_i)^2} \right) - y_2 = 0, \end{cases} \tag{3}$$

where a_i and b_i are the coordinates of the radius-vector \vec{l}_i i.e. $\vec{l}_i = (a_i, b_i)$.

From an algebraic point of view, system (3) is a system of two rational equations (over a field of real numbers) from two unknowns, which are given in Cartesian coordinates on the R^2 plane. The system (3) will be considered, just above the field of complex numbers C , while we denote it by (3a). The set of solutions of system (3) is obviously the set of real solutions of system (3a). We note that all the coefficients of the equations of system (3a) are real.

In terms of algebraic geometry, the image of a source in an N -point gravitational lens can be defined as follows:

Definition. An image of a point source in an N -point gravitational lens will be called the real solution of system (3a) without regard to multiplicity. The set of images is the set of different real solutions of this system.

3. REDUCTION OF THE PROBLEM TO THE FUNDAMENTAL PROBLEM OF CLASSICAL ALGEBRAIC GEOMETRY

The main problem of classical algebraic geometry is the problem of studying systems of polynomial equations.

Let us investigate the set of solutions of system (3a). To do this, we transform the equations of the system to a polynomial form

$$\begin{cases} F_1 = (x_1 - y_1) \prod_{i=1}^N h_i - \sum_{j=1}^N m_j (x_1 - a_j) \prod_{i=1, i \neq j}^N h_i = 0, \\ F_2 = (x_2 - y_2) \prod_{i=1}^N h_i - \sum_{i=1}^N m_i (x_2 - b_i) \prod_{i=1, i \neq j}^N h_i = 0, \end{cases} \quad (4)$$

where $h_i = (x_1 - a_i)^2 + (x_2 - b_i)^2, i = 1, 2, \dots, N$.

The polynomial form of the equations in system (4) is necessary for its investigation by methods of algebraic geometry.

We shall consider the equations of the system (4) over the field C of complex numbers in the affine coordinate system C^2 . The system (4) is not equivalent to the system (3a), but it follows from it. The set of solutions of the system (3a) can be obtained from the set of solutions of the system (4). For this, by removing solutions from it in which the system (3a) is not defined. These solutions are pairs of numbers that are the coordinates of the point masses. Indeed, we directly verify that the points with coordinates $(a_i, b_j), i = 1, \dots, N$, is a solution of the system (4), but the system (3a), in these points is not defined.

Let f_1 and f_2 be the left-hand sides of the first and second equations of system (3a), $M(f_1, f_2)$ be the solution set of system (3a), $V(F_1, F_2)$ be the solution set of system (4), and $\text{Re } V(F_1, F_2) \subset V(F_1, F_2)$ the subset of its real solutions, then we have

$$M(f_1, f_2) = \text{Re } V(F_1, F_2) / \{\cup(a_i, b_i)\}. \quad (5)$$

From the theorem on the structure of the set of solutions of a system of polynomial equations, see [11] it follows that the set $V(F_1, F_2)$ can be represented in the form

$$V(F_1, F_2) = (V^0(F_1, F_2)) \cup (V^1(F_1, F_2)), \quad (6)$$

where $V^1(F_1, F_2)$ is the set of solutions depending on a single parameter, and $V^0(F_1, F_2)$ is the discrete set of solutions of system (3a).

The set $V^0(F_1, F_2)$ is obviously discrete and, moreover, finite. The sets have dimension $\dim V^0(F_1, F_2) = 0$ and $\dim V^1(F_1, F_2) = 1$.

4. STUDY OF THE SET $V^1(F_1, F_2)$ (EXTENDED SOLUTIONS)

A number of theorems, which allow us to determine if the set $V^1(F_1, F_2)$ is empty, see, for example, [12, 13]. In [5], we give an algorithm that allows us to describe this set analytically, if it is not empty. If the set $V^1(F_1, F_2)$ is not empty, then the equations of system (3) are said to have a common component. The equation of the common component can be obtained from the analytical description of the set $V^1(F_1, F_2)$.

4.1. 1-POINT LENS (SCHWARZSCHILD LENS)

We apply the theorems presented in the Appendix for constructing the set $V^1(F_1, F_2)$, in the case of a single-point gravitational lens.

Let the 1-point lens have coordinates $a_1 = 0, b_1 = 0$. Let $L: R_Y^2 \rightarrow R_X^2$ be the transformation from the plane of the source to the plane of the lens, determined by the system of equations

$$\begin{cases} y_1 = x_1 - \frac{x_1}{x_1^2 + x_2^2}, \\ y_2 = x_2 - \frac{x_2}{x_1^2 + x_2^2}. \end{cases} \quad (7)$$

Equations of the system are defined for all points such that $x_1^2 + x_2^2 \neq 0$, that is, except for the origin of the point $O(0, 0)$. At the origin, the inverse mapping is not defined. But if we transform system (7) to polynomial form

$$\begin{cases} (x_1^2 + x_2^2)(x_1 - y_1) - x_1 = 0, \\ (x_1^2 + x_2^2)(x_2 - y_2) - x_2 = 0, \end{cases} \quad (8)$$

then the inverse transformation of $L^{-1}: R_X^2 \rightarrow R_Y^2$ is completely determined by the system of equations

$$\begin{cases} x_1^3 + x_1 x_2^2 - x_1^2 y_1 - x_2^2 y_1 - x_1 = 0, \\ x_1^2 x_2 + x_2^3 - x_1^2 y_2 - x_2^2 y_2 - x_2 = 0. \end{cases} \quad (9)$$

We calculate the result R_1 by the variable x_1 for which we represent the equation in lexicographic form

$$\begin{cases} x_1^3 - y_1 x_1^2 + (x_2^2 - 1)x_1 - x_2^2 y_1 = 0, \\ (x_2 - y_2)x_1^2 + x_2^3 - x_2^2 y_2 - x_2 = 0. \end{cases} \quad (10)$$

Result by degree x_1 has the form

$$R_1 = \begin{vmatrix} 1 & r_{12} & r_{13} & r_{14} & 0 \\ 0 & 1 & r_{12} & r_{13} & r_{14} \\ r_{21} & 0 & r_{23} & 0 & 0 \\ 0 & r_{21} & 0 & r_{23} & 0 \\ 0 & 0 & r_{21} & 0 & r_{23} \end{vmatrix}, \quad (11)$$

where $r_{12} = -y_1, r_{13} = x_2^2 - 1, r_{14} = -x_2^2 y_1, r_{21} = x_2 - y_2, r_{23} = x_2^3 - x_2^2 y_2 - x_2$.

We have

$$R_1 = -x_1^3 y_1^2 + x_2^2 y_1^2 y_2 + x_2 y_1^2 - x_2^3 y_2^2 + x_2^2 y_2^3. \quad (12)$$

In order for the system equations (7) to have a common component, we need $R_1 \equiv 0$.

Applying Theorem 5a for the decomposition of R_1 on indecomposable components, we have

$$x_2(-(y_1^2 + y_2^2)x_2^2 + y_2(y_1^2 + y_2^2)x_2 + y_1^2) \equiv 0. \quad (13)$$

The equation is divided into two equations

$$\begin{aligned} x_2 &\equiv 0, \\ -(y_1^2 + y_2^2)x_2^2 + y_2(y_1^2 + y_2^2)x_2 + y_1^2 &\equiv 0. \end{aligned} \quad (14)$$

Each of the equations is considered to be a polynomial of the variable x_2 .

A polynomial is identically equal to zero if and only if all its coefficients are equal to zero.

From here there is a system of equations

$$\begin{cases} y_1^2 + y_2^2 = 0, \\ y_1^2 = 0. \end{cases} \quad (15)$$

Next we have that $y_1 = 0, y_2 = 0$. Substituting in (7) we have

$$\begin{cases} x_1 - \frac{x_1}{x_1^2 + x_2^2} = 0, \\ x_2 - \frac{x_2}{x_1^2 + x_2^2} = 0 \end{cases} \Rightarrow \begin{cases} x_1 \left(1 - \frac{1}{x_1^2 + x_2^2}\right) = 0, \\ x_2 \left(1 - \frac{1}{x_1^2 + x_2^2}\right) = 0. \end{cases} \quad (16)$$

The system (16) decomposes into three systems and one equation:

$$\begin{cases} x_1 = 0, \\ x_2 = 0, \end{cases} \quad (17a)$$

$$\begin{cases} x_1 = 0, \\ 1 - \frac{1}{x_1^2 + x_2^2} = 0, \end{cases} \quad (17b)$$

$$\begin{cases} 1 - \frac{1}{x_1^2 + x_2^2} = 0, \\ x_2 = 0, \end{cases} \quad (17c)$$

$$1 - \frac{1}{x_1^2 + x_2^2} = 0. \quad (18)$$


System (17a) has a solution $x_1 = 0, x_2 = 0$ but this solution is not a solution of system (7), since the system equations at the point $O(0, 0)$ are not defined.

The system (17b) has two solutions $x_1 = 0, x_2 = \pm 1$.

The system (17c) has two solutions $x_1 = \pm 1, x_2 = 0$.

We have the transformation of equation (18) into

$$x_1^2 + x_2^2 - 1 = 0. \quad (19)$$

Equation (19) is the equation of an  individual circle in the plane X with a center at the point $O(0, 0)$.

The solution of systems (17b) and (17c) satisfies equation (19).

The solution of system (8) is the coordinates of the points of single circle with center at the point $O(0, 0)$.

Equation (19) is the equation of the general component, hence the set

$$V^1(F_1, F_2) = \{x_1, x_2 \mid x_1^2 + x_2^2 - 1 = 0\}. \quad (20)$$

In the same way, we compute the resultant R_2 . By virtue of the symmetry of variables, we have the same solution.

4.2. 2-POINT LENS

We research a two-point gravitational lens with equal masses $m_1 = m_2 = \frac{1}{2}$.

The masses are on the abscissa at a distance a from the origin of coordinates.

In this case, system (3) looks like this:

$$\begin{cases} y_1 = x_1 - \frac{1}{2} \frac{x_1 - a}{(x_1 - a)^2 + x_2^2} - \frac{1}{2} \frac{x_1 + a}{(x_1 + a)^2 + x_2^2}, \\ y_2 = x_2 - \frac{1}{2} \frac{x_2}{(x_1 - a)^2 + x_2^2} - \frac{1}{2} \frac{x_2}{(x_1 + a)^2 + x_2^2}. \end{cases} \quad (21)$$

We transform the equation of system (21) into a polynomial form, and represent the obtained polynomials F_1 and F_2 in lexicographic form with increasing degrees of variable x_1 :

$$\begin{cases} F_1 = -a^2(a^2 + 2x_2^2 y_1) \\ \quad + (a^2 - x_2^2 + (a^2 + x_2^2))x_1 + 2y_1(a^2 - x_2^2)x_1^2 \\ \quad + (2x_2^2 - 1 - 2a^2)x_1^3 + y_1 x_1^4 + x_1^5, \\ F_2 = (-a^2 x_2 - x_2^3 + (a^2 + x_2^2)^2(x_2 - y_2)) \\ \quad - (x_2 + 2(a^2 - x_2^2)(x_2 - y_2))x_1^2 + (x_2 - y_2)x_1^4, \end{cases} \quad (22)$$

We will remove from the system the variable x_1 , using the resultant $R_1 = R(F_1, F_2)$.

Sylvester matrix $S_1 = S(F_1, F_2)$ has order (for x_1) $\deg F_1 + \deg F_2 = 9$.

Because $R_1 = \det S_1$, we have

$$\begin{aligned} R_1 = & 4a^4 x_2^2 (a^2 + x_2^2) (-a^2 y_2^3 \\ & + (a^2 y_2^2 - y_2^2 - 4a^4 y_2^2 - 4a^2 y_2^4) x_2 \\ & + (-4a^2 y_2 + 4a^4 y_2 - 4a^6 y_2 - y_1^2 y_2^2 \\ & - 4a^2 y_1^2 y_2 + 8a^4 y_1^2 y_2 - 4a^2 y_1^2 y_2 \\ & - 5y_2^3 + 4a^2 y_2^3 - 8a^4 y_2^3 - 4a^2 y_2^5) x_2^2 \\ & + (-4a^4 + 4a^6 + y_1^2 + 4a^2 y_1^2 \\ & - 8a^4 y_1^2 + 4a^2 y_1^4 + y_2^2 - 12a^2 y_2^2 \\ & + 8a^4 y_2^2 - 8y_1^2 y_2^2 - 8y_2^4 + 4a^2 y_2^4) x_2^3 \\ & - 4(a^4 y_2 - a^2 y_2 - y_1^2 y_2 - 2a^2 y_1^2 y_2 \\ & + y_1^4 y_2 - y_2^3 + 2a^2 y_2^3 + 2y_1^2 y_2^3 + y_2^5) x_2^4 \\ & + 4(a^4 - 2a^2 y_1^2 + y_1^4 + 2a^2 y_2^2 + 2y_1^2 y_2^2 + y_2^4) x_2^5. \end{aligned} \quad (23)$$

In order for the system equation to have a common component, it is sufficient that the objects $R_1 \equiv 0$.

We have that the equation decomposes into three simple equations and one non-trivial equation.

From the trivial equation $a^4 \equiv 0, x_2^2 \equiv 0, (a^2 + x_2^2) \equiv 0$ it follows that their solutions are reduced to 1-lens, or incommensurate.

We have a nontrivial equation

$$\begin{aligned} & -a^2y_2^3 + (a^2y_2^2 - y_2^2 - 4a^4y_2^2 - 4a^2y_2^4)x_2 \\ & + (-4a^2y_2 + 4a^4y_2 - 4a^6y_2 - y_1^2y_2^2 - 4a^2y_1^2y_2 + 8a^4y_1^2y_2 \\ & \quad - 4a^2y_1^2y_2 - 5y_2^3 + 4a^2y_2^3 - 8a^4y_2^3 - 4a^2y_2^5)x_2^2 \\ & + (-4a^4 + 4a^6 + y_1^2 + 4a^2y_1^2 - 8a^4y_1^2 + 4a^2y_1^4 + y_2^2 \\ & \quad - 12a^2y_2^2 + 8a^4y_2^2 - 8y_1^2y_2^2 - 8y_2^4 + 4a^2y_2^4)x_2^3 \\ & - 4(a^4y_2 - a^2y_2 - y_1^2y_2 - 2a^2y_1^2y_2 + y_1^4y_2 \\ & \quad - y_2^3 + 2a^2y_2^3 + 2y_1^2y_2^3 + y_2^5)x_2^4 \\ & + 4(a^4 - 2a^2y_1^2 + y_1^4 + 2a^2y_2^2 + 2y_1^2y_2^2 + y_2^4)x_2^5 = 0. \end{aligned} \tag{24}$$

We equate all coefficients to zero, and have a system of equations

$$\left\{ \begin{aligned} & -a^2y_2^3 = 0, \\ & a^2y_2^2 - y_2^2 - 4a^4y_2^2 - 4a^2y_2^4 = 0, \\ & -4a^2y_2 + 4a^4y_2 - 4a^6y_2 - y_1^2y_2^2 \\ & \quad - 4a^2y_1^2y_2 - 4a^2y_2^5 - 4a^2y_1^2y_2 \\ & \quad - 5y_2^3 + 4a^2y_2^3 - 8a^4y_2^3 + 8a^4y_1^2y_2 = 0, \\ & -4a^4 + 4a^6 + y_1^2 + 4a^2y_1^2 - 8a^4y_1^2 + 4a^2y_1^4 + y_2^2 \\ & \quad - 12a^2y_2^2 + 8a^4y_2^2 - 8y_1^2y_2^2 - 8y_2^4 + 4a^2y_2^4 = 0, \\ & a^4y_2 - a^2y_2 - y_1^2y_2 - 2a^2y_1^2y_2 + y_1^4y_2 \\ & \quad - y_2^3 + 2a^2y_2^3 + 2y_1^2y_2^3 + y_2^5 = 0, \\ & a^4 - 2a^2y_1^2 + y_1^4 + 2a^2y_2^2 + 2y_1^2y_2^2 + y_2^4 = 0. \end{aligned} \right. \tag{25}$$

We have

$$\left\{ \begin{aligned} & a = 0, \\ & y_2 = 0, \\ & -y_1^2y_2^2 - 5y_2^3 = 0, \\ & -y_1^2y_2 + y_1^4y_2 - y_2^3 + 2y_1^2y_2^3 + y_2^5 = 0, \\ & y_1^4 + 2y_1^2y_2^2 + y_2^4 = 0. \end{aligned} \right. \tag{26}$$

The system has one solution $a = 0, y_1 = 0, y_2 = 0$. Hence this solution reduces the 2-point gravity lens to 1-point.

Similarly we calculate the resultant R_2 as

$$\begin{aligned} R_2 = & 4a^4(a - x_1)x_1^2(a + x_1)(-a^2y_1^3 + \\ & + (1 + a^2 + 4a^4 - 4a^2y_1^2 - 4a^2y_2^2)y_1^2x_1 \\ & + (-4a^2 + 4a^4 - 4a^6 + 5y_1^2 + 4a^2y_1^2 + 8a^4y_1^2 - 4a^2y_1^4 \\ & + y_2^2 - 4a^2y_2^2 - 8a^4y_2^2 - 8a^2y_1^2y_2^2 - 4a^2y_2^4)y_1x_1^2 \\ & + (4a^4 + 4a^6 + (-1 - 12a^2 - 8a^4 + 8y_1^2 + 4a^2y_1^2)y_1^2 \end{aligned}$$

$$\begin{aligned} & + (-1 + 4a^2 + 8a^4 + 8y_1^2 + 8a^2y_1^2 + 4a^2y_2^2)y_2^2)x_1^3 \\ & + (4a^2y_1^2 + a^4y_1 - y_1^3 - 2a^2y_1^3 + y_1^5 - y_1y_2^2 \\ & \quad + 2a^2y_1y_2^2 + 2y_1^3y_2^2 + y_1y_2^4)x_1^4 \\ & - 4(a^4 - 2a^2y_1^2 + y_1^4 + 2a^2y_2^2 + 2y_1^2y_2^2 + y_2^4)x_1^5). \end{aligned} \tag{27}$$

We have that the solution of system (21) reduces the 2-point gravity lens to 1-point. Whence

- for the 1-point gravitational lens set we have

$$V^1(F_1, F_2) = \{ x_1, x_2 \mid x_1^2 + x_2^2 - 1 = 0 \};$$

- for the 2-point gravitational lens set we have

$$V^1(F_1, F_2) = \emptyset.$$

Based on the studies we have carried out above, one can prove that there are no extended objects for N -point gravitational lenses, i.e. $V^1(F_1, F_2) = \emptyset$.

The set $M(f_1, f_2)$ can be represented in the form

$$M(f_1, f_2) = M^0(f_1, f_2) \cup M^1(f_1, f_2), \tag{28}$$

where $M^0(f_1, f_2) = \text{Re } V^0(F_1, F_2) / \{\cup(a_i, b_i)\}$ and $M^1(f_1, f_2) = \text{Re } V^1(F_1, F_2) / \{\cup(a_i, b_i)\}$.

It is known that the set $M^1(f_1, f_2)$, for a point source in 1-point lens is not empty, see for example [9, 10, 14], coincides with $V^1(F_1, F_2)$, see [5] and is Einstein ring. But for a point source in symmetric 2-point lens, we proved [5] that the set $M^1(f_1, f_2)$ is empty and put forward hypothesis: for N -point lens this set is empty for $N > 1$.

5. THE STUDY OF THE SET $V^0(F_1, F_2)$ (POINT SOLUTIONS)

To research the set of solutions $V^0(F_1, F_2)$ of system (3) we use the Bezout theorem, see for example [11–13, 15].

In most monographs, the authors formulate the Bezout theorem in geometric terms; see for example [11, 12, 15]. One of these theorems is quoted in Appendix.

In [13] Bezout's theorem is formulated in algebraic terms, but for equations given in affine coordinates. This theorem is also quoted in Appendix.

For our purposes, we formulate this theorem in algebraic terms, but for functions given in homogeneous coordinates.

Theorem 1 (Bezout). *Let $G_1(X_0 : X_1 : X_2)$ and $G_2(X_0 : X_1 : X_2)$ be homogeneous polynomials, $\deg G_1(X_0 : X_1 : X_2) = n, \deg G_2(X_0 : X_1 : X_2) = m$ and the resultant $R_1(G_1, G_2)$, with respect to variable X_1 not identically equal to zero. Then the resultant $R_1(G_1, G_2)$ is a homogeneous polynomial with respect to variables X_0 and X_2 , and $\deg R_1(G_1, G_2) = n \cdot m$.*

Proof. The resultant $R_1(G_1, G_2)$ is a polynomial in the variables X_0 and X_2 . We denote it by F , and write

$F = R_1(G_1, G_2)$. Let us prove that the polynomial $F = F(X_0, X_2)$ is homogeneous and of degree $\deg F = n \cdot m$. Really, we have

$$F(tX_0, tX_2) = R_1(G_1(tX_0 : X_1 : tX_2), G_2(tX_0 : X_1 : tX_2)). \quad (29)$$

According to Theorem 1a, see Appendix, we have

$$R_1(G_1(tX_0 : X_1 : tX_2), G_2(tX_0 : X_1 : tX_2)) = \det \text{Sul}(G_1(tX_0 : X_1 : tX_2), G_2(tX_0 : X_1 : tX_2)). \quad (30)$$

where the right-hand side of expression (30) is the determinant of the Sylvester matrix. The order of this determinant is $n + m$.

Elements of the Sylvester matrix are coeficies from the lexicographic representation of homogeneous polynomials G_1 and G_2 . We have that

$$G_1 = \sum_{i=0}^n a_i X_1^{n-i} \quad \text{and} \quad G_2 = \sum_{j=0}^m b_j X_1^{m-j}.$$

The coefficients $a_i = a_i(X_0 : X_2)$ and $b_j = b_j(X_0 : X_2)$ are homogeneous polynomials.

The degree is $\deg a_i(X_0 : X_2) = i$ and $\deg b_j(X_0 : X_2) = j$, that is, $a_i(tX_0 : tX_2) = t^i a_i(X_0 : X_2)$ and $b_j(tX_0 : tX_2) = t^j b_j(X_0 : X_2)$.

We multiply every row of the determinant of the Sylvester matrix by the parameter t in some degree. We choose a degree so that all elements of the column are of the same degree with respect to t .

We multiply the i -th row of the determinant of the Sylvester matrix by t^i where $i = 1, 2, \dots, m$, and j -th row by t^j where $j = m + 1, m + 2, \dots, m + n$. We take the factor t^s from the s -th column where $s = 1, 2, \dots, m + n$. We denote the total power of t by S . We have

$$S = \sum_{i=1}^{m+n} i - \sum_{i=1}^n i - \sum_{i=1}^m i = \frac{(n+m)(n+m+1)}{2} - \frac{n(n+1)}{2} - \frac{m(m+1)}{2} = nm. \quad (31)$$

In this way,

$$\det \text{Sul}(G_1(tX_0 : X_1 : tX_2), G_2(tX_0 : X_1 : tX_2)) = t^{nm} \det \text{Sul}(G_1(X_0 : X_1 : X_2), G_2(X_0 : X_1 : X_2)). \quad (32)$$

The determinant of the Sylvester matrix does not depend on the parameter t .

Substituting (32) in (30) and further in (29) we have

$$F(tX_0, tX_2) = t^{nm} F(X_0, X_2). \quad (33)$$

Consequently, the resultant is a homogeneous polynomial of degree nm . ■

Theorem 1 admits a generalization. We have proved an analogous assertion for systems of equations of several variables, see [7, 16].

We transform system (3) and apply Bezout's theorem to its study.

In the equations of system (3) we proceed to homogeneous coordinates. Let

$$\begin{cases} x_1 = X_1/X_0, \\ x_2 = X_2/X_0. \end{cases} \quad (34)$$

After reducing the equations of the system to a polynomial form, we have

$$\begin{cases} X_0^{2N+1} F_1\left(\frac{X_1}{X_0}, \frac{X_2}{X_0}, y_1\right) = \Phi_1(X_0 : X_1 : X_2) = 0, \\ X_0^{2N+1} F_2\left(\frac{X_1}{X_0}, \frac{X_2}{X_0}, y_2\right) = \Phi_2(X_0 : X_1 : X_2) = 0. \end{cases} \quad (35)$$

The coordinates X_0, X_1, X_2 are obviously projective coordinates.

The system (34) defines surjective mapping, $\mathfrak{S}: C^2 \rightarrow CP^2$. The triple of complex numbers $(X_0 : X_1 : X_2)$ are the coordinates of the point and defines the point $p \in CP^2$ in the projective plane CP^2 . The triple $(\lambda X_0 : \lambda X_1 : \lambda X_2)$ specifies the same point if $\lambda \neq 0$.

Therefore we have the following result.

Theorem 2. *The system of polynomial equations*

$$\begin{cases} \Phi_1(X_0 : X_1 : X_2) = 0, \\ \Phi_2(X_0 : X_1 : X_2) = 0 \end{cases} \quad (36)$$

has in the projective plane CP^2 , counting multiplicity, exactly $m \cdot n$ solutions, where, $m = \deg \Phi_1$, and $n = \deg \Phi_2$, if $\gcd(\Phi_1, \Phi_2)$ belongs to the coefficient field C .

Functions $\Phi_1 = \Phi_1(X_0 : X_1 : X_2)$ and $\Phi_2 = \Phi_2(X_0 : X_1 : X_2)$ are homogeneous functions of degree $2N + 1$.

If, at least one of the coordinates of the point p is equal to zero, say that this point is irregular. Otherwise, the point is called regular.

A straight line that consists of irregular points is called an irregular line.

The projective plane CP^2 has three irregular straight lines, which are given by the equations

$$X_0 = 0, \quad X_1 = 0, \quad X_2 = 0. \quad (37)$$

The set of points CP^2 one of the coordinates, which is equal to the number $h \neq 0$, is called affine map on CP^2 and denoted by $A^2(h)$. The complement of a $CP^2 \setminus A^2(h)$ consists of a one-dimensional complex projective subspace, which is called an infinitely distant line of the affine map, see for example [11, 15]. The infinitely distant line of any affine map $A^2(h)$ is evidently irregular.

In particular, if we put $X_0 = 1$, then the set of points CP^2 with coordinates $(1 : X_1 : X_2)$ will be

affine map of $A^2(1)$, and the infinity of the straight line of this map will be given by equation $X_0 = 0$.

Consider the situation of general position, i.e. the source is not on the caustic. In this case, the Jacobian of the system of lens equations is not equal to zero.

Theorem 3. *In a situation of general position (the Jacobian of the system of lens equations is not equal to zero), the number of point images in an N -point gravitational lens has parity opposite to the parity of the number N .*

In the proof of Theorem 3 we use the following lemma.

Lemma 1. *The number of irregular solutions of system (37), on line $X_0 = 0$, is $2N$.*

Proof. Using (37), we reduce the system to the form

$$\left\{ \begin{array}{l} (X_1 - X_0 y_1) \prod_{i=1}^N H_i \\ - X_0^2 \sum_{j=1}^N m_j (X_1 - X_0 a_j) \prod_{i=1, i \neq j}^N H_i = 0, \\ (X_2 - X_0 y_2) \prod_{i=1}^N H_i \\ - X_0^2 \sum_{j=1}^N m_j (X_2 - X_0 b_j) \prod_{i=1, i \neq j}^N H_i = 0, \end{array} \right. \quad (38)$$

where $H_i = (X_1 - X_0 a_i)^2 + (X_2 - X_0 b_i)^2$.

Let $X_0 = 0$. We have

$$\left\{ \begin{array}{l} X_1 \prod_{i=1}^N (X_1^2 + X_2^2) = 0, \\ X_2 \prod_{i=1}^N (X_1^2 + X_2^2) = 0 \end{array} \right. \Rightarrow \left\{ \begin{array}{l} X_1 (X_1^2 + X_2^2)^N = 0, \\ X_2 (X_1^2 + X_2^2)^N = 0 \end{array} \right. \Rightarrow (X_1^2 + X_2^2)^N = 0$$

$$\Rightarrow X_1 = \pm i X_2 \Rightarrow \begin{cases} X_1 = c, \\ X_2 = \pm ic. \end{cases} \quad (39)$$

Finally we have two N -fold solutions: $P_1 = (0 : a : ic)$ and $P_2 = (0 : a : -ic)$. ■

Proof of Theorem 2. For the degrees of the polynomials of systems (3) and (5) we have $\deg F_1 = \deg F_2 = \deg \Phi_1 = \deg \Phi_2 = 2N + 1$.

By Bezout's theorem, the system of equations (36) has $(2N + 1)^2$ solutions, which include an even number of $2q$ complex conjugate solutions and $P = 2N$ irregular solutions.

Therefore, the number of real solutions of system (36),

$$\begin{aligned} \text{card}(\text{real } V^0(F_1, F_2)) &= (2N + 1)^2 - 2q - P \\ &= (2N + 1)^2 - 2q - 2N \\ &= 4N^2 + 2N + 1 - 2q. \end{aligned} \quad (40)$$

From the fact that the restriction of the inverse mapping $\mathfrak{S}^{-1}: CP^2 \rightarrow C^2$ to the affine map $A^2(1)$ is a bijection that is given by the equations $X_0 = 1, X_1 = x_1, X_2 = x_2$, we have

$$\begin{aligned} \text{card}(M^0(f_1, f_2)) &= \text{card}(\text{real } V^0(F_1, F_2)) - N \\ &= 4N^2 + N + 1 - 2q. \end{aligned} \quad (41)$$

In a situation of general position, the point source is not on the caustic, therefore, all elements of the set $\text{real } V(f_1, f_2)$ are different.

In this case, each point of the set $\text{real } V(f_1, f_2)$ is, by definition, an image.

It follows from (9) that the parity of the number of images is opposite to the parity of the number N . ■

Theorem 3 does not contradict the theorem on the oddness of the number of images in transparent lenses [9, 10].

Example 1. For a 1-point lens, the number of images is 2, see [9, 10, 14].

Example 2. For a 2-point lens, the number of images is 3 or 5 see [17].

6. CONCLUSIONS

Applying methods of algebraic geometry, we constructed an algorithm that separates images of dimensions 1 and 0.

In the present paper, for an image of dimension 1, it is proved that for single-point sources there exists a unique image of dimension 1-the Einstein ring; Einstein's ring is only in a single-point lens; the point source in other lenses does not have images of dimension 1 for $N > 1$. For an image of dimension 0, it is proved that in any N -point gravitational lens: there are a finite number of images; the parity of the number of images is always the opposite of the parity of the number N .

The assertion for the number of images of dimension 0 was proved by us earlier, see [18]. In [18], we used the geometric method of algebraic geometry-the Newton diagram. In the present paper all the assertions are proved algebraically. This opens the possibility, to use N -point gravitational lenses, not only approximate or numerical methods, but also computer algebra systems.

A. APPENDIX

Let $f(x, y)$ be a function of two variables, and $f(x, y)$, at the point (x_0, y_0) , n -times continuous, differentiable function. Then Taylor's formula holds:

$$f(x, y) = f(x_0, y_0) + \sum_{k=1}^n f^{(k)}(x - x_0, y - y_0) + r_n(x, y), \quad (42)$$

Theorem 4a. Let $f(x_1, x_2) = \sum_{i,j=0}^{i+j \leq n} f_{ij} x_1^i x_2^j$ and $g(x_1, x_2) = \sum_{i,j=0}^{i+j \leq m} g_{ij} x_1^i x_2^j$ be polynomials. Let their coefficients be such that $f_{n0} \neq 0$, $f_{0n} \neq 0$, $g_{m0} \neq 0$, $g_{0m} \neq 0$. Then

$$\begin{aligned} \deg R_1(f(x_1, x_2), g(x_1, x_2)) \\ = \deg f(x_1, x_2) \deg g(x_1, x_2) = nm. \end{aligned} \quad (49)$$

For the proof see, e.g., [12].

Theorem 5a. The polynomials $F_1(x_1, x_2)$, $F_2(x_1, x_2)$ have a non-trivial common component if and only if $R_1(F_1(x_1, x_2), F_1(x_1, x_2))$ or $R_2(F_1(x_1, x_2), F_1(x_1, x_2))$.

For the proof see, e.g., [12].

Definition 5a. A formal sum $G = G(x_1, x_2)$ of the form $G = \sum_{i,j=0}^{i+j \leq n} g_{ij} x_1^i x_2^j$ is called a polynomial n -form of variables x_1, x_2 over the field K . that is, G is a polynomial of degree n in variables x_1, x_2 with indefinite coefficients g_{ij} in the field K .

The expression “the function will be sought in the form of a n -form” is usually understood as a procedure for determining the undetermined coefficients of a given n -form.

Theorem 6a (Criterion of non-decomposability). Let F be a polynomial in the variables over the field of complex numbers and G . Let R_1 and R_2 be the resultant of the variables x_1 and x_2 over the field of complex numbers. Let R_1 and R_2 be the resultants with respect to the variables x_1 and x_2 respectively.

The polynomial F is not decomposable if and only if the systems of equations

$$\begin{aligned} \frac{\partial^i}{\partial x_1^i} R_1(F(x_1, x_2), G(x_1, x_2)) &= 0, \\ \frac{\partial^i}{\partial x_2^i} R_2(F(x_1, x_2), G(x_1, x_2)) &= 0, \end{aligned} \quad (50)$$

where $i = 1, \dots, m$, $m = \deg R_1(F, G)$, have only zero solutions.

The system (50) is considered with respect to the undetermined coefficients g_{ij} n -form G , as with respect to unknown variables.

For the proof see, e.g., [7].

The theorem admits a generalization to the case of systems of equations of several variables see [7, 16].

The criterion was formulated and proved by the authors earlier [7].

REFERENCES

- [1] E. Y. Bannikova, A. T. Kotvytskiy. Three Einstein rings: explicit solution and numerical simulation. *MNRAS* **445**(4):4435–4442, 2014. DOI:10.1093/mnras/stu2068.
- [2] A. T. Kotvytskiy. Gravitational lensing by straight cosmic strings. *TMP* **184**(1):160–174, 2015. DOI:10.1007/s11232-015-0315-x.
- [3] A. Cassan. An alternative parameterisation for binary-lens caustic-crossing events. *Astronomy and Astrophysics* **491**(2):587–595, 2008. DOI:10.1051/0004-6361:200809795.
- [4] P. Schneider, A. Weiss. The two-point-mass lens: detailed investigation of a special asymmetric gravitational lens. *Astronomy and Astrophysics* **164**(2):237–259, 1986.
- [5] A. T. Kotvytskiy, S. D. Bronza, K. D. Nerushenko, V. Y. Shablenko. Matematychnii zmist kiltisia Einsteina ta umovy yogo vynyknennja. doslidzhennja uzahal'nyh umov[in ukrainian]. In *Zbirnyk naukovykh prac' VI-i Mizhrehional'noi nauково-praktychnoi konferencii "Astronomiya i s'ogodennya"*, pp. 198–213. Vinnytsia, 2017.
- [6] A. T. Kotvytskiy, S. D. Bronza, S. R. Vovk. Estimating the number of images N-point gravitational lenses algebraic geometry methods. *Visnyk KhNU, serija "Fizyka"* **24**:55–59, 2016.
- [7] S. D. Bronza, A. T. Kotvytskiy. Mathematical bases of the theory of N-point gravitational lenses. Part 1. Elements of algebraic geometry. *Visnyk KhNU, serija "Fizyka"* **26**:6–32, 2017.
- [8] A. T. Kotvytskiy, S. D. Bronza. Mathematical bases of the theory of N-point lenses. *Odessa Astronomical Publications* **29**:31–33, 2016. DOI:0.18524/1810-4215.2016.29.84958.
- [9] A. F. Zakharov. *Gravitacionnye linzy i mikrolinzy [in Russian]*. Janus-K, Moscow, 1997.
- [10] P. Schneider, J. Ehlers, E. Falco. *Gravitational lenses*. Second printing. Springer-Verlag Berlin Heidelberg, 1999.
- [11] I. V. Arzhantcev. *Bazisy Grebnera i sistemy algebraicheskikh uravnenij [in Russian]*. MCNMO, Moscow, 2003.
- [12] R. J. Walker. *Algebraic curves*. Springer-Verlag New York, 1978.
- [13] E. Kalinina, A. J. Uteshev. *Teoria iskluchenij [in Russian]*. NII khimii SPbGU Saint Petersburg, 2002.
- [14] P. V. Bliokh, A. A. Minakov. *Gravitational Lenses [in Russian]*. Naukova Dumka, Kiev, 1989.
- [15] M. Reid. *Undergraduate Algebraic Geometry*. Math Inst., University of Warwick, 2013.
- [16] S. D. Bronza. Kriteriy neprivodimosti mnogochlenov ot dvuh peremennykh nad polem kompleksnih chisel[in russian]. In *Zbirnyk naukovykh prac'*, pp. 114–115. UkrDUZT, 2016.
- [17] S. H. Phie. Infimum microlensing amplification of the maximum number of images of N-point lens systems. *The astrophysical journal* **484**(1):63–69, 1997. DOI:10.1086/304336.
- [18] A. T. Kotvytskiy, S. D. Bronza, V. Y. Shablenko. Correlation of the number of images of an N-point gravitational lens and the number of solutions of its system. In *Astronomy and beyond: astrophysics, cosmology, cosmomicrophysics, astroparticle physics, radioastronomy and astrobiology*, p. 12. Odessa International Astronomical Gamow Conference-School, 2017.
- [19] B. L. van der Waerden. *Modern Algebra I, II*. New York:Frederic Ungar Publishing Co., 1950.
- [20] S. Lang. *Algebra*. Revised Third Edition. Columbia University. New York, 1965.